This article was downloaded by:

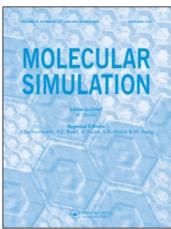
On: 14 January 2011

Access details: Access Details: Free Access

Publisher Taylor & Francis

Informa Ltd Registered in England and Wales Registered Number: 1072954 Registered office: Mortimer House, 37-

41 Mortimer Street, London W1T 3JH, UK



Molecular Simulation

Publication details, including instructions for authors and subscription information: http://www.informaworld.com/smpp/title~content=t713644482

Crystal structure prediction could have helped the experimentalists with polymorphism in benzamide

Juergen Thun^a; Markus Schoeffel^a; Josef Breu^a

^a Department of Inorganic Chemistry I, University of Bayreuth, Bayreuth, Germany

To cite this Article Thun, Juergen , Schoeffel, Markus and Breu, Josef (2008) 'Crystal structure prediction could have helped the experimentalists with polymorphism in benzamide', Molecular Simulation, 34:10,1359-1370

To link to this Article: DOI: 10.1080/08927020802208943 URL: http://dx.doi.org/10.1080/08927020802208943

PLEASE SCROLL DOWN FOR ARTICLE

Full terms and conditions of use: http://www.informaworld.com/terms-and-conditions-of-access.pdf

This article may be used for research, teaching and private study purposes. Any substantial or systematic reproduction, re-distribution, re-selling, loan or sub-licensing, systematic supply or distribution in any form to anyone is expressly forbidden.

The publisher does not give any warranty express or implied or make any representation that the contents will be complete or accurate or up to date. The accuracy of any instructions, formulae and drug doses should be independently verified with primary sources. The publisher shall not be liable for any loss, actions, claims, proceedings, demand or costs or damages whatsoever or howsoever caused arising directly or indirectly in connection with or arising out of the use of this material.



Crystal structure prediction could have helped the experimentalists with polymorphism in benzamide

Juergen Thun, Markus Schoeffel and Josef Breu*

Department of Inorganic Chemistry I, University of Bayreuth, Bayreuth, Germany (Received 30 January 2008; final version received 16 May 2008)

Benzamide was the first molecular material for which polymorphism was reported, as long as 176 years ago. Unfortunately, due to very similar cell metrics leading to massive peak overlap, the metastable form reported by Liebig escaped structural characterisation by XRD until recently. With the help of crystal structure prediction this old riddle of 'Liebig's' polymorph of benzamide could have been solved many years earlier. Performing state of the art crystal structure prediction runs applying a commercial suite of programs (Cerius² and the MS Modeling) identified both the metastable and the thermodynamically stable phase reported by Liebig as low energy packings.

Keywords: crystal structure prediction; density functional theory; point charges; benzamide, polymorphism

1. Introduction

Polymorphism is a widespread phenomenon occurring with many known crystalline materials. Polymorphism describes the ability of crystalline material to crystallise in more than one crystal structure [1]. The physical and chemical properties like melting point, solubility and bioavailability may differ. The latter feature is of course of great interest in pharmaceutical research. As a consequence of these differing properties the crystal structure is relevant for patent licensing as well. The most famous and often cited example in this respect is the case of Ritonavir [2,3]. In most cases, an active pharmaceutical ingredient (API) is protected by a patent covering both, molecular and crystal structure(s). If additional crystal structures/ metastable phases of this API can be obtained, this patent protection can be circumvented. Even a separate patent covering this new crystal structure can be applied for. In the dye industry the situation is comparable [4,5]. For these reasons a good knowledge about possible polymorphs is essential. Crystal structure prediction (CSP) represents a fast and cheap method of acquiring a general idea of local energy minima on the lattice energy hyper surface. Several different techniques can be found in the literature and the success of these methods is checked regularly [6,7,8]. Rigid body calculations allow a very accurate description of the intermolecular potentials using distributed multipole analysis [9,10,11]. But these techniques are more costly and at the same time a view on the results presented here would suggest that the precision of the force field cannot be improved significantly. Things become even more complicated when flexible bodies are used during the minimisation

procedure. The use of quantum mechanically derived monopole charges is much less demanding and yields comparable accuracies in respect to lattice parameters, as will be shown later on. The energy ranking will be more exact using distributed multipoles [12] but for a first overview of possible polymorphs and their packing motifs monopole charges are adequate. Nevertheless, CSP still is a challenging and tricky task to perform, especially for the API molecules the pharmaceutical industry is interested in. But even small molecules may still pose challenges. Benzamide appears to be a very good benchmark for CSP because it is a small molecule with only one torsional degree of freedom (see Figure 1). In total, by now three polymorphs of this small molecule have been identified experimentally [13]. The lattice energies of two phases reported by Liebig (phase I and phase III) lie at close quarters and the plate like crystals of phase I have been regarded as to be the global minimum structure. Phase II reported by Davey et al. [16] appears to be a highly metastable or even transient packing. The task was to reproduce these known polymorphs by CSP. A state of the art approach was taken using different force fields (point charge models) in combinations with the commercial Cerius² [14] and MS Modeling suite of programs [15] for global optimisation. Of course, all predictions had been performed without implementing any knowledge about the different packing motifs found experimentally for these three polymorphs.

2. Brief summary of crystallographic details

All three polymorphs (see Table 1) show a similar hydrogen bonding motif [16,17,13]. Two benzamide molecules form

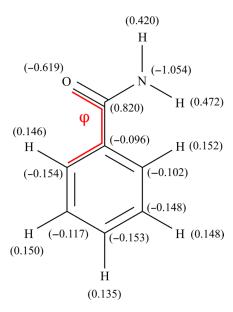


Figure 1. A benzamide molecule showing the torsion angle ϕ and the Gaussian03 calculated partial charges in brackets.

hydrogen-bonded dimers and are connected by further hydrogen bonds to form long ribbons along the short axis of the unit cell (see Figure 2). Interestingly, despite the similar hydrogen bonding motifs, the $\pi-\pi$ interactions are quite different. In phase I a so-called shifted π stack is realised, while phase II and phase III show a herringbone pattern like arrangement (see Figure 3). So you can find all low energy $\pi-\pi$ interaction motifs realised within the different crystal structures of the same molecule. One major difference lies in the number of symmetry independent molecules. While phase I and phase III contain only one independent molecule in the asymmetric unit cell, phase II contains two independent molecules. CSP runs with more than one independent molecule still represent a challenge for the established algorithms [6].

3. Computational methods

3.1 Method of generating the molecular structure and charge

3.1.1 Geometry

During the global optimisation applying a Monte-Carlo (MC) algorithm as implemented in the Cerius² [14] and

Table 1. Crystallographic details of phase I [17], phase II [17] and phase III [13].

	Phase I	Phase II	Phase III
a/Å	5.607	17.4317(2)	5.0551(1)
b/Å	5.046	14.1944(3)	5.5141(1)
c/Å	22.053	4.977(1)	22.9565(9)
β / $^{\circ}$	90.66	_	101.292(2)
Volume/Å ³	624.25	1231.4(3)	627.50(3)
Space group	$P2_1/c$	Pba2	$P2_1/c$
Z^{\prime}	1	2	1
Z	4	8	4

MS Modeling suite of programs [15], all molecules are held rigid. Therefore, in CSP runs for molecules that possess internal degrees of freedom, several different starting conformers should be used to get a reliable sampling over the whole potential energy surface. Since benzamide has only one torsional degree of freedom, low energy starting conformers could easily be identified by a grid search technique. The torsion angle ϕ (see Figure 1) was modified from 0° up to 90° and a geometry optimisation of the gas phase molecules was performed with the Forcite module within the Material Studio® software [15] using Gasteiger charges [18] and the Dreiding 2.21 force field [19]. Unsurprisingly, only one minimum structure could be found at 0° torsion. As a second starting conformer, the saddle point on the gas phase energy hyper surface was chosen which has 90° torsion. Please note, that these molecular structures used in the MC runs differ significantly in respect to the torsion angles from the experimentally found ones (about 30°). No bias regarding the molecular structure found experimentally was allowed in the CSP runs.

3.1.2 Choice of force field

Different density functional theory (DFT) methods had been used to derive several sets of partial charges that best represent the molecular electrostatic potential of benzamide. In Material Studio DMol³ [20,21] was used with a DNP basis set and the GGA-PBE functional and also the GGA-PW91 functional. The atomic charges were then fitted so as to reproduce the electrostatic potential (ESP)

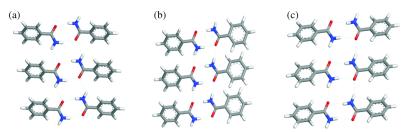


Figure 2. Hydrogen bonded networks. View along the short axis of the unit cells. (a) phase I; (b) phase II and (c) phase III. Dashed blue lines show the hydrogen bonds.

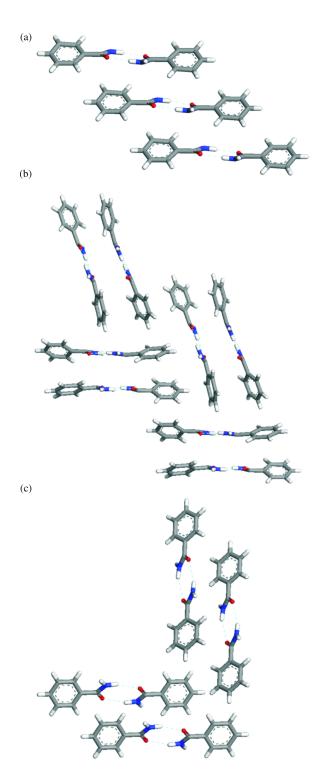


Figure 3. A projection along the ribbon axes shows the $\pi-\pi$ interactions in the different forms. (a) phase I, view along [010]; (b) phase II, view along [001] and (c) phase III, view along [100].

[22,23]. Additionally, ESP derived partial charges have been obtained in the same manner by using Gaussian03 [24] with a 6-311++G(d,p) basis set and the PBE1PBE[25] or the B3LYP hybrid functional [26,27], respectively. Both Gaussian03 calculations produced the

same set of partial charges. Next, a quality check of the calculated sets of point charges in combination with different force fields had been performed using the crystal structure of the accurately known phase I as a benchmark. The molecules had been treated as flexible bodies during the geometry optimisation within the Forcite module and no further constraints had been introduced. The best results as judged by the deviation in cell constants could be obtained by using the Dreiding 2.21 force field in combination with the Gaussian03 point charges. Using the COMPASS force field or partial charges based on DMol³ calculations lead to larger deviations. The details can be seen in Table 2. When comparing the experimental lattice parameters with the parameters obtained after lattice energy minimisations deviations of less than 5% were regarded acceptable. The good agreement achieved using the Gaussian03 set of partial charges in combination with the Dreiding 2.21 force is reassuring and thus this force field was used for all subsequent calculations. Special caution has been taken to assure that the cost function in both program suites applied in global optimisation was indeed identical. For instance, as Ewald summation was used throughout all geometry optimisations and lattice energy calculations for both Coulomb and dispersion forces, it was assured that the same set of accuracies (0.000001 kcal/mol) was applied.

3.2 Generating trial crystal structures

3.2.1 Cerius²

Monte-Carlo simulated annealing algorithm [28,29,30,31] as implemented in the Cerius² software suite was used. The nine most commonly observed space groups (P2₁/c, P1̄, P2₁2₁2₁, P2₁, C2/c, Pbca, Pnma, Pna2₁ and Pbcn) were searched with one molecule in the asymmetric unit. The space groups $P\bar{1}$, $P2_1$, $P2_1/c$, $P2_12_12_1$ and $Pna2_1$ were also searched with Z'=2. The space group Pba2, the space group experimentally observed for phase II, is not available for CSP in the used Cerius² suite of programs. Therefore, additionally a search in Pna2₁ was performed. Pna2₁ is a subgroup of Pba2 and can be represented by a unit cell doubled along the c-axis. Consequently, this further search in Pna2₁ had to be performed with Z'=4. Rigid bodies were used throughout the simulated annealing procedure. Two different conformations with a 0° or 90° torsion angle ϕ , respectively, were used as rigid bodies. Because the balance between intra- and intermolecular energies was neglected during the MC runs, initially rather loosely packed trial crystal structures are generated. At least three independent repeats had been performed with each of the two different starting conformers until a convergence in low energy structures could be found. Each run consists of 7000 MC steps and yielded a acceptance ratio of about 60%.

Table 2. Relative energies and unit cell parameters^a of the observed and energy minimised crystal structures of phase I using Gaussian03 PBE1PBE ESP-charges (Gaus) and DMol³ PBE ESP-charges (DMol) in combination with the Dreiding 2.21 (drei) and COMPASS (com) force field within MS Modeling.

	a/Å	b/Å	c/Å	β/°	ρ /g cm ⁻³	Energy ^b /kcal mol ⁻¹
Obs.	5.607	5.046	22.053	90.66	1.287	
Gaus _{drei}	-2.2	-2.3	4.4	0.7	1.293	-63.859
$DMol_{drei}$	6.6	1.8	-2.1	1.7	1.285	-48.388
Gaus _{com}	- 8.6	14.7	- 1.2	- 11.3	1.261	-80.207
$\mathrm{DMol}_{\mathrm{com}}$	- 3.8	2.1	0.0	- 5.3	1.223	-33.048

^aExperimental lattice parameters and, in italics, lattice parameters for minimised structures given as a percentage difference from the experimental values. ^bEnergy per molecule.

3.2.2 MS Modeling

All above mentioned space groups had been searched as well with the Monte-Carlo algorithm as implemented in the MS Modeling suite of programs. Additionally, the space group Pba2 had been examined by using Z'=2. In all space groups at least three independent repeats had been performed with each of the two different starting conformers until a convergence in low energy structures was assured. As in the Cerius² calculations, 7000 MC steps had been performed and also yielded about 60% of acceptance.

3.2.3 Geometry optimisation

For both, the Cerius² and the MS Modeling MC algorithms, a clustering algorithm limits the number of accepted crystal structures to 500 after each simulated annealing run. The geometry optimisation of 500 structures can be handled with ordinary cpu-power within a few hours or days depending on the degrees of freedom in the explored space group. The clustering of the structures is based on the similarity of their radial distribution functions. The molecular flexibility was introduced directly after the clustering and all accepted crystal structures were energy minimised with the force field described before. At this stage, all structural parameters like lattice parameters, molecular positions, orientations and conformations were relaxed. After geometry optimisation, the density of all crystal structures reaches 1.1–1.4 g cm⁻¹ compared to 0.3–0.8 g cm⁻¹ before the minimisation.

3.2.4 Ranking of structures

After a final clustering, the calculated crystal structures were ranked on the one hand according to their total energy and on the other hand according to their density. At least 20 structures with low energy and 20 structures with high density were also compared 'by eye' attempting to identify the packing motifs of the experimentally known polymorphs. The unit cells of interesting predicted crystal structures were transformed into the standard settings applying LePage as implemented in Platon [32].

4. Results and discussion

4.1 Prediction for $\mathbf{Z}' = 1$

4.1.1 Cerius² results

The predicted structures with the lowest lattice energies were found in the space group $P2_1/c$ and two structures clearly break away in the energy density plot (see Figure 4) with very similar energy and density. These two predicted structures are the most densely packed crystal structures and have the lowest achieved energy that have been found in all runs and they can be assigned to phase I and phase III, respectively (see Table 3). The details of the predicted and the experimentally observed packing motifs, a shifted π stack for phase I and a herringbone pattern for phase III, are in close agreement (see Figure 5). The maximum error of the predicted lattice parameters amounts to 5.7% and also the experimentally observed torsion angles ϕ were predicted very well within errors of about 8 and 5% for phase I and III, respectively. Please note that the starting conformers used during MC exhibited either a 0° or 90° torsion angle ϕ and the point charges had been calculated and fitted to the electrostatic potential of the gas phase minimum. Taking this into account, the agreement of predicted and experimental torsion angles is astonishing. It is expected that a more accurate description of the

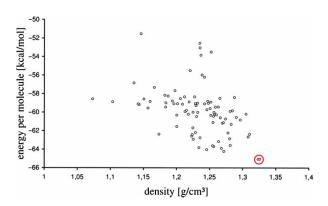


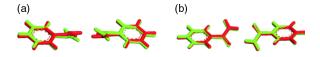
Figure 4. Energy-density plot in the space group $P2_1/c$ (Cerius² runs). Packings corresponding to phase I and phase III are circled.

		Phase I			Phase III	
Parameter	Cerius ²	Exp.	Δ/%	Cerius ²	Exp.	Δ/%
a/Å	5.318	5.607	-5.2	4.902	5.055	-3.0
b/Å	4.905	5.046	-2.8	5.232	5.514	-5.1
b/Å c/Å	23.315	22.053	5.7	23.778	22.957	3.6
β / $^{\circ}$	91.71	90.66	1.2	95.75	101.29	-5.5
$\rho/g \text{ cm}^{-3}$	1.323	1.287	2.8	1.326	1.282	3.4
$\phi/^{\circ}$	30.54	33.20	-8.0	30.18	28.84	4.6

Table 3. Comparison of the predicted and observed structural parameters and total energies for phase I and phase III (Cerius² calculation).

Energy^a/kcal mol⁻¹

^aEnergy per molecule.



-65.132

Figure 5. Overlay of the Cerius² predicted (red) and observed (green) dimers of (a) phase I and (b) phase III.

electrostatic interactions could be achieved using distributed multipole moments. However, results recently published by Karamertzanis et al. show only slightly smaller deviations [33]. Even with our very simple charge model, the density and even torsions could be reproduced within nearly the same error limits. For a small molecule like benzamide, there seems to be no need for a more accurate description of the electrostatics, which of course would be much more time consuming. Even the energy ranking of both predicted structures is in agreement with experiment (see Table 3). The lowest energies of predicted structures in the other searched space groups had been around $-62 \, \text{kcal mol}^{-1}$ per molecule and the density was less than $1.28 \, \text{g cm}^{-3}$.

4.1.2 MS Modeling results

-65.055

The experimentally observed polymorphs, phase I and phase III, could be identified as low energy/high density structures with the MS Modeling suite of programmes as well (see Figure 6). The predicted crystal structures compare very well with the experimentally observed ones (see Table 4 and Figure 7) and structural differences between the Cerius² predicted structures and the MS Modeling predicted ones are very small (compare Tables 3 and 4).

Also, a final single point energy calculation applying Cerius² on phase I and phase III as predicted in MS Modeling yielded within calculational errors the same lattice energies as the Cerius² predicted ones (compare Table 3 and Table 5) and the gradients were negligible. However, the absolute values between Cerius² and MS Modeling differ by roughly 1 kcal/mol as can be seen in Table 5. The difference occurs within the intramolecular part of the total energy. Although we intended to apply the very same force field parameters and precision in calculations, surprisingly the intramolecular fraction of the total energy is more positive within MS Modeling.

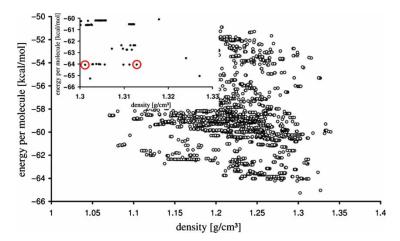


Figure 6. Energy-density plot in the space group $P2_1/c$. (MS Modeling runs). The inset shows a close up of the high density and low energy region. Packings corresponding to phase I and phase III are circled.

Table 4. Comparison of the predicted and observed structural parameters and total energies for phase I and phase III (MS Modeling 4.0 calculation).

		Phase I			Phase III			
Parameter	MS Mod.	Exp.	2/%	MS Mod.	Exp.	7/√	lowest energy structure	highest density structure
a/Å	5.322	5.607	-5.1	4.904	5.055	-3.0	3.477	11.787
b/Å	4.905	5.046	-2.8	5.232	5.514	-5.1	7.795	5.914
$c/ iny{ m \AA}$	23.316	22.053	5.7	23.776	22.957	3.6	18.912	12.777
βl°	91.66	99.06	1.1	95.83	101.29	- 5.4	145.06	137.37
ρ /g cm ⁻³	1.301	1.287	1.1	1.326	1.282	3.4	1.302	1.334
<i>مار</i>	30.72	33.20	-7.5	30.17	28.84	4.6	13.49	21.27
Energy ^a /kcal mol ⁻¹	-64.065			-64.052			-65.070	-59.858
		1				1		

^aEnergy per molecule.

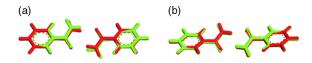


Figure 7. Overlay of the MS Modeling predicted (red) and observed (green) dimers of (a) phase I and (b) phase III.

Unfortunately, in the Version 4.0 of MS Modeling there is no force field editor for the Dreiding force field, which rendered a final analysis, of the origin of this divergence, impossible. This difference in lattice energies comes down to a different balance/weighting of interand intramolecular energies in Cerius² and MS Modeling. This in turn might explain, why the global optimisation runs identified different local minima (see below).

There is also a quite pronounced difference in the energy ranking of the generated crystal structures. While in the Cerius² prediction phase I and III could be found on position one and two, respectively, the energy ranking position after the MS Modeling prediction has been 3 and 12, respectively. The insert of Figure 6 shows that many structures with a similar energy had been predicted, but the arrangement of the molecules is quite different. A density ranking placed both predicted phase I and phase III on position 30 for phase I and 13 for phase III, respectively. This result is not influenced by the molecular structure used in the prediction. The starting conformer with a 90° torsion angle φ yielded similar results. Taking a closer look on the predicted lowest energy structure pronounced differences in the packing motif as compared with the known phases are obvious, especially in respect to the hydrogen bonding pattern (see Figure 8(a)). This structure contains no dimers, the benzamide molecules only form hydrogenbonded chains along the short axis of the unit cell. This catamer synthon is frequently found for carbonic acids [34]. The contribution of the hydrogen bonding to the electrostatic energy is less attractive than in the above mentioned crystal structures of phase I and III, but the van der Waals interaction is slightly more favourable leading to a lower total energy. This suggests that the odds to crystallise this polymorph should increase in strongly hydrogen bonding solvents that are capable to break up the benzamide dimers in solution.

Interestingly, the most densely packed structure also contains chains of hydrogen-bonded molecules along the short axis of the unit cell, but in this case some kind of a zig-zag orientation was realised (see Figure 8(b)). The hydrogen bonding part of the non-bonding energy again is less attractive as compared to the energies of the experimentally known phase I and phase III, because the strong hydrogen bonds within the benzamide dimers are not realised. Looking on the $\pi-\pi$ interactions, a shifted π stack had been realised for the lowest energy structure, while the

	Pha	se I	Phase III	e III
Energy ^a kcal mol ⁻¹	MS Mod	Cerius ²	MS Mod.	Cerius ²
Intramoleculare	8.478	7.411	8.451	7.448
Bond	1.974	1.776	1.982	1.753
Angle	2.948	2.547	2.945	2.565
Torsion	3.428	2.958	3.389	2.993
Inversion	0.128	0.130	0.135	0.137
Intermoleculare	-72.543	-72.543	-72.503	-72.503
Hydrogen bond	-6.744	-6.744	-6.735	-6.735
Van der Waals	2.855	2.855	2.821	2.821
Electrostatic	-68.654	-68.654	-68.589	-68.589
Total	-64.065	-65.132	-64.052	-65.055

Table 5. Single point intra- and intermolecular energies of phase I and phase III as predicted in MS Modeling.

highest density structure shows a tilted herringbone pattern arrangement. One C—H bond is pointing to the middle of a neighbouring benzene ring. Nevertheless, the total energy of this structure is about 5 kcal mol⁻¹ per molecule higher compared to the lowest energy structure.

4.1.3 Discussion

Although the methods should be nearly identical, the results of the calculations with the two programs appear quite diverse and might lead to different conclusions on either the choice of method or more generally on the usefulness of crystal structure prediction. On a first glance, Cerius² appears to give a much clearer picture: The known phase I and phase III are found as the lowest energy crystal structures – a result that would be very useful if it was general and reproducible also for other molecules. The same set of calculations performed within the MS

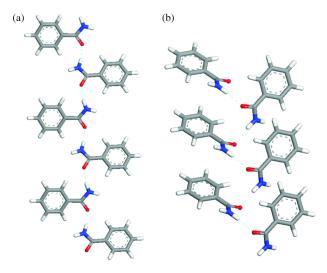


Figure 8. Hydrogen bonding pattern in the predicted structures with (a) the lowest energy and (b) the highest density.

Modeling suite of programs also finds phase I and III, but quite far up the list of structures – a less useful result? Apparently, the global optimisation of MS Modeling is more thorough and is capable of identifying additional low energy/high density packings. As mentioned above, a possible reason for the better performance of the global search might be the differing balance/weighting of intra-and intermolecular contributions to the total energy.

A final single point energy calculation of the MS Modeling predicted lowest energy structure within the Cerius² suite shows clearly, that the global optimisation within Cerius² seems to be incomplete. This structure is assigned the lowest energy also within Cerius² $(-66.053\,\mathrm{kcal\ mol}^{-1}$ per molecule) and hence Cerius² failed to identify this valid global minimum.

In summary, two alternative conclusions may be drawn:

On one hand, the experimental screening may not be exhausted. This would be the more expedient result for experimentalists and might possibly lead to the identification of another polymorph. This result would pass the ball back to the experimentalists and ask for further polymorph screening experiments.

On the other hand, if the experimental screening has indeed been exhausted, it would suggest that the cost function applied is not sufficiently precise and therefore the ranking of phase I and III is wrong. This would pass the ball into the field of the prediction community and asks for further improvements of force fields.

Nevertheless, could this excellent agreement of the predicted and the experimentally observed crystal structures have helped to earlier identify the metastable polymorph reported by Liebig? This phase III managed to 'hide' behind the stable phase I for such a long time because the metric of the unit cell is very similar to the one of phase I and peak overlap is severe [13]. Since phase III can only be obtained as a 50:50 mixture with phase I, it was therefore simply overlooked.

^aEnergy per molecule.

Obviously, even such small deviations in cell axes and relative coordinates as observed in this CSP greatly affects both, the positions and the intensities of the diffraction peaks. The differences between the calculated and experimentally observed X-ray powder diffraction pattern (PXRD) are immense (Figure 9). It is hard to imagine that any person or algorithm would have been able to identify phase III in the overlapping experimental PXRD trace of the mixture accessible in the experiment with the help of the PXRD calculated using the predicted structure. The lattice energy as calculated using classical force fields is inappropriate to be used as the only cost function for structure solution in real space. However, it is obvious from the calculated diffraction patterns that the very first low angle peaks (around $5^{\circ} 2\theta$) of the two predicted phases are overlapping while in the experimental trace a shoulder is visible. Consequently, with some trust in the CSP results it could have appeared that the experimental trace might not be pure phase I as has been assumed for such a long time but a mixture of the two low energy structures.

4.2 Prediction for $\mathbb{Z}' > 1$

4.2.1 Cerius² results

The experimentally observed crystal structure of phase II could not been found in either of the prediction runs in $Pna2_I$ with Z'=4 using starting conformers with 0° or 90° torsion angle. Therefore, starting conformers with torsions as observed in the known packing motif of phase II had been included as a bias in the MC runs, again without success. Additionally, benzamide dimers and pairs of dimers as observed in the crystal structure of phase II had

also been applied as rigid units during the Monte-Carlo algorithm. The ESP-derived partial charges had been calculated for these oligomeric units with Gaussian03 and the afore-mentioned basis set and functionals. Despite all these additional efforts that have been severely biased by knowledge from the experimentally determined structure, phase II could not been obtained by CSP.

A search in the space group $P\bar{1}$ with Z'=2 found phase I and phase III. The predicted crystal structures fit quite well with experimental structures, but the prediction of phase III showed bigger deviations as compared to the runs in space group $P2_1/c$ and the energies were higher (see Table 6). These calculations showed, that a crystal structure prediction with more than one molecule within the asymmetric unit is in principle possible for benzamide.

4.2.2 MS Modeling results

Although in MS Modeling, prediction runs in space group Pba2 and Z'=2 were accessible, phase II could not be identified either. The same biased rigid units as described in the preceding paragraph had been used, again without any mentionable results.

However, as with Cerius², phase I and phase III can also be identified in a search in the space group $P\bar{1}$ and two independent molecules within the asymmetric unit (see Table 7). This suggests that the problems encountered in predicting phase II are not necessarily related with the fact that one has to deal with Z'=2. Preliminary results on quenched MD simulation runs starting with the experimentally known structure did not yield additional insights.

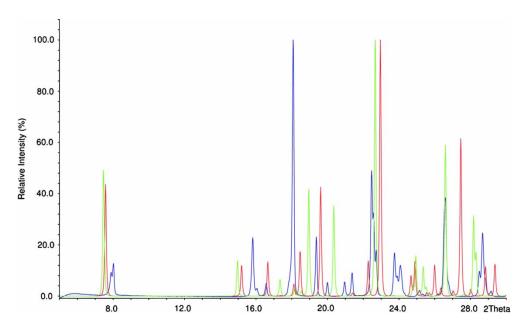


Figure 9. Experimental powder diffraction pattern (blue) combined with the calculated diffraction data of the predicted phase I (red) and phase III (green).

Table 6. Comparison of the predicted and observed structural parameters and total energies for phase II and phase III in $P\bar{1}$ (Cerius² calculation).

		Phase I			Phase III	
Parameter	Cerius ²	Exp.	Δ/%	Cerius ²	Exp.	Δ/%
a/Å	5.559	5.607	-0.0	4.921	5.055	-2.7
b/Å	4.923	5.046	-2.4	5.207	5.514	-5.6
b/Å c/Å	23.132	22.053	4.9	25.044	22.957	9.1
β / $^{\circ}$	91.64	90.66	1.0	93.98	101.29	-7.7
$\rho/g \text{ cm}^{-3}$	1.272	1.287	-1.2	1.257	1.282	-2.0
$\phi/^{\circ}$	30.76	33.20	-7.3	29.91	28.84	3.7
Energy ^a /kcal mol ⁻¹	-64.223			-63.348		

^aEnergy per molecule.

4.2.3 A closer look at phase II

To shed some light into the difficulties encountered with the prediction of phase II a simple static energy optimisation of the known phase II had been performed. The crystal structure of phase II had to be determined with a diffraction trace that suffered from two major draw backs: A large and broad background caused by the solvent and peak overlap with the majority phase I (82.7% [17]). Therefore, it may be anticipated that the details of the crystal structure are less reliable. And indeed, the experimental crystal structure represents a saddle point in respect to the lattice energy as defined by the force field. Especially the van der Waals energy is highly positive (+66.75 kcal mol⁻¹ per molecule). Apparently there are some non-bonding contacts that are too short leading to highly repulsive forces that direct the static minimisation into a wrong secondary minimum. To avoid this, a stepwise approach was used to relax the experimentally known crystal structure of phase II. In the first step, all C—H bonds had been shortened to 0.9 Å and additionally the lattice parameters of the a- and b-axes had been inflated by about 10%. In the following two rigid body optimisations, no internal energies had been considered. During the first geometry optimisation only the orientation

of the rigid bodies had been allowed to relax within a fixed unit cell and during the second geometry optimisation additionally the lattice parameters had been optimised. A third optimisation run relaxed the molecules within the fixed unit cell and finally all parameters had been optimised. The final energy minimised crystal structure is close to the experimental structure. However, the deviation in the b axis still amounts to 10.5%. This suggests that the dimers cannot be accurately represented by standard force fields and this may contribute to the difficulties encountered in CSP of phase II. Apparently, phase II is very sensitive to small changes in the molecular geometry and therefore CSP is more severely hampered by inaccuracies in the force field and the negligence of balance of the inter-/intramolecular interactions during the MC runs. The calculated energy of phase II is very similar to the energies of phase I and phase III (see Table 8) and also the energy ranking is in good agreement with experimental results. Please note, that the energies published recently had been calculated without taking the internal energies into account (compare Thun et al. [13]). An overlay of the optimised and the experimental crystal structure shows the good agreement after the minimisation. In the relaxed structure, the hydrogen bonding within

Table 7. Comparison of the predicted and observed structural parameters and total energies for phase I and phase III in $P\bar{1}$ (MS Modeling calculation).

	Phase I			Phase III		
Parameter	MS Mod	Exp.	Δ/%	MS Mod	Exp.	Δ/%
a/Å	5.329	5.607	-5.0	4.922	5.055	-2.6
b/Å c/Å	4.907	5.046	-2.8	5.451	5.514	-1.1
c/Å	23.256	22.053	5.5	23.730	22.957	3.4
β / $^{\circ}$	91.74	90.66	1.2	91.72	101.29	-9.4
$\rho/g \text{ cm}^{-3}$	1.324	1.287	2.9	1.313	1.282	2.4
φ/°	33.15	33.20	0.0	32.98	28.84	14.4
Energy ^a /kcal mol ⁻¹	-64.052			-63.335		

^aEnergy per molecule.

Table 8. Cell parameters of the observed and energy minimised crystal structure of phase II using Gaussian03 PBE1PBE ESP-charges, the Dreiding 2.21 force field and a stepwise approach.

Parameter	Cerius ²	Exp.	Δ /%
a/Å	17.095	17.432	- 1.9
b/Å	15.684	14.194	10.5
c/Å	5.151	4.977	3.5
$\rho/g \text{ cm}^{-3}$	1.165	1.307	-10.9
$\phi_1/^\circ$	24.68	38.77	36.3
$\phi_2/^\circ$	27.58	13.72	101.0
Energy ^a /kcal mol ⁻¹	-64.355	116.677 ^b	

^a Energy per molecule. ^b Single point energy calculation in Cerius² prior to lattice energy minimisation.

a benzamide dimer is more distinct as compared to the experimental structures due to a more favourable orientation of the two-amide groups (see Figure 10 and compare the change in the torsion angles ϕ_1 and ϕ_2 of both independent molecules in Table 8). In Figure 11 the herringbone motif is shown before and after refinement. After minimisation, the π -systems of two dimers lie perfectly above each other and a C-H group of one benzene ring points to the centre of a neighbouring benzene system. Nevertheless, also this relaxed structure could not been found in either of the performed simulations. Since the agreement in lattice constants of experimental and optimised phase II is mediocre, expectedly a comparison of the simulated and experimental powder pattern also shows pronounced differences (see Figure 12).

But clearly, phase II is a valid local minimum in the energy landscape of benzamide packings. In respect to energy

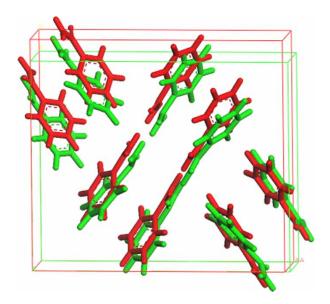


Figure 10. Overlay of the optimised (red) and experimental (green) unit cell of phase II.

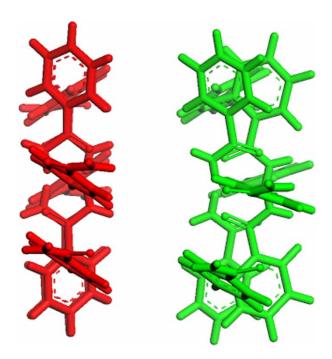


Figure 11. View perpendicular onto a pair of dimers of benzamide as found in the crystal structure of phase II before (green) and after (red) energy minimisation.

it is close enough to phase I and phase III in order to be identified in the CSP runs. At the moment it is hard to say whether the problems with CSP of this known polymorph are related to the uncertainties in the force field or the difficulties encountered in global optimisation with Z' = 2.

5. Conclusion

This paper presents the successful prediction of two benzamide polymorphs. Using crystal structure prediction,

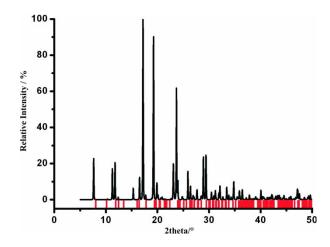


Figure 12. Simulated powder pattern of the minimised structure and ticks (red) of the experimentally known phase II.

the metastable polymorph described by Wöhler and Liebig 176 years ago could have been solved many years earlier. Within one day of calculation the structure of this metastable phase III would have been predicted within the Cerius² suite of programs. Applying the Material Studio Modeling suite of programs, phase I and phase III could be identified as well. However, many more packings with comparable energy and/or density are predicted rendering the identification of the experimentally known structures tricky. The lowest energy structure found with MS Modeling is not yet known for benzamide but represents a well known structure motif for carbonic acids. Additional experimental polymorph screenings will be performed to further investigate the experimentally accessible phase space.

Lattice energy minimisations suggest that phase II is very sensitive to small changes in the molecular geometry. At the moment it is hard to say whether the inability to identify phase II with CSP is related to the uncertainties in the force field or the difficulties encountered in global optimisations with Z' = 2.

During all presented simulations, ESP-derived atom centred partial charges had been used and the chosen representation of the electrostatic potential seemed to work well for all performed rigid and flexible body minimisations. Even the experimentally observed torsion angles ϕ could be reproduced quite accurate. Due to the fact, that three different $\pi-\pi$ interactions can be found within the crystal structures of one molecule, benzamide is a good benchmark system for further investigations on the influence of different force fields and charge distributions. Finally it can be said, a good choice of monopole charges together with different starting conformers provides a fast and efficient way to get an overview over possible crystal structures.

Acknowledgements

We would like to thank the anonymous referee for his/her keen interest in the work, evident from the indepth and most helpful comments which helped to improve the manuscript further indeed.

References

- J. Bernstein, Polymorphism in Molecular Crystals, Oxford University Press, Oxford, 2002.
- [2] S.R. Chemburkar, J. Bauer, K. Deming, H. Spiwek, K. Patel, J. Morris, R. Henry, S. Spanton, W. Dziki, W. Porter et al., *Dealing with the impact of ritonavir polymorphs on the late stages of bulk drug process development*, Org. Process Res. Dev. 4 (2000), pp. 413–417.
- [3] H. Nowell and S. L. Price, Validation of a search technique for crystal structure prediction of flexible molecules by application to piracetam, Acta Crystallogr. B 61 (2005), pp. 558–568.
- [4] M.U. Schmidt, C. Buchsbaum, J.M. Schnorr, D.W.M. Hofmann, and M. Ermrich, *Pigment-Orange-5: crystal structure determination* from a non-indexed X-ray powder diagram, Z. Kristallogr. 222 (2007), pp. 30–33.

- [5] M.U. Schmidt, R.E. Dinnebier, and H. Kalkhof, Crystal engineering on industrial diaryl pigments using lattice energy minimizations and X-ray powder diffraction, J. Phys. Chem. B 111 (2007), pp. 9722–9732.
- [6] G.M. Day, W.D.S. Motherwell, H.L. Ammon, S.X.M. Boerrigter, R.G. Della Valle, E. Venuti, A. Dzyabchenko, J.D. Dunitz, B. Schweizer, B.P. van Eijck et al., A third blind test of crystal structure prediction, Acta Crystallogr. B 61 (2005), pp. 511–527.
- [7] J.P.M. Lommerse, W.D.S. Motherwell, H.L. Ammon, J.D. Dunitz, A. Gavezzotti, D.W.M. Hofmann, F.J.J. Leusen, W.T.M. Mooij, S.L. Price, B. Schweizer et al., A test of crystal structure prediction of small organic molecules, Acta Crystallogr. B 56 (2000), pp. 697–714.
- [8] W.D.S. Motherwell, H.L. Ammon, J.D. Dunitz, A. Dzyabchenko, P. Erk, A. Gavezzotti, D.W.M. Hofmann, F.J.J. Leusen, J.P.M. Lommerse, W.T.M. Mooij et al., Crystal structure prediction of small organic molecules: a second blind test, Acta Crystallogr. B 58 (2002), pp. 647–661.
- [9] U. Koch and A.J. Stone, Conformational dependence of the molecular charge distribution and its influence on intermolecular interactions, J. Chem. Soc. Faraday Trans. 92 (1996), pp. 1701–1708.
- [10] A.J. Stone and M. Alderton, Distributed multipole analysis methods and applications, (Reprinted from Molecular Physics, vol 56, pg 1047–1064, 1985) Mol. Phys. 100 (2002), pp. 221–233.
- [11] A.J. Stone, Distributed multipole analysis or how to describe a molecular charge-distribution, Chem. Phys. Lett. 83 (1981), pp. 233–239.
- [12] G.M. Day, W.D.S. Motherwell, and W. Jones, A strategy for predicting the crystal structures of flexible molecules: the polymorphism of phenobarbital, Phys. Chem. Chem. Phys. 9 (2007), pp. 1693–1704.
- [13] J. Thun, L. Seyfarth, J. Senker, R.E. Dinnebier, and J. Breu, Polymorphism in benzamide: Solving a 175-year-old riddle, Angew. Chem. Int. Ed. Engl. 46 (2007), pp. 6729–6731.
- [14] Accerlys Cerius², Release 4.9, Accelrys Software, Inc, San Diego, 2003.
- [15] Accerlys Material Studio, Release 4.0, Accelrys Software, Inc, San Diego. 2006.
- [16] N. Blagden, R. Davey, G. Dent, M. Song, W.I.F. David, C.R. Pulham, and K. Shankland, Woehler and Liebig revisited: A small molecule reveals its secrets – the crystal structure of the unstable polymorph of benzamide solved after 173 years, Cryst. Growth Des. 5 (2005), pp. 2218–2224.
- [17] W.I.F. David, K. Shankland, C.R. Pulham, N. Blagden, R.J. Davey, and M. Song, *Polymorphism in benzamide*, Angew. Chem. Int. Ed. 44 (2005), pp. 7032–7035.
- [18] J. Gasteiger and M. Marsili, *Iterative partial equalization of orbital electronegativity a rapid access to atomic charges*, Tetrahedron 36 (1980), pp. 3219–3228.
- [19] S.L. Mayo, B.D. Olafson, and W. A. Goddard, *Dreiding a generic force-field for molecular simulations*, J. Phys. Chem. 94 (1990), pp. 8897–8909.
- [20] B. Delley, An all-electron numerical-method for solving the local density functional for polyatomic-molecules, J.Chem.Phys. 92 (1990), pp. 508–517.
- [21] G. Fitzgerald, B. Delley and J.W. Andzelm, DFT: From molecules to solids using LCAO numerical orbitals, Abstracts Of Papers Of The American Chemical Society 220 (2000), p. U187
- [22] B.H. Besler, K.M. Merz, and P.A. Kollman, Atomic charges derived from semiempirical methods, J. Comput. Chem. 11 (1990), pp. 431–439.
- [23] U.C. Singh and P.A. Kollman, *An approach to computing electrostatic charges for molecules*, J. Comput. Chem. 5 (1984), pp. 129–145.
- [24] M.J. Frisch, G.W. Trucks, H.B. Schlegel, G.E. Scuseria, M.A. Robb, J.R. Cheeseman, J.A. Montgomery, Jr, T. Vreven, K.N. Kudin, J.C. Burant et al. *Gaussian 03, Revision C.02*, Gaussian, Inc, Wallingford CT, 2004.
- [25] J.P. Perdew, K. Burke, and M. Ernzerhof, Generalized gradient approximation made simple, Phys. Rev. Lett. 77 (1996), pp. 3865–3868.

- [26] A.D. Becke, *Density-functional thermochemistry.3*. The role of exact exchange, J. Chem. Phys. 98 (1993), pp. 5648–5652.
- [27] A.D. Becke, A new mixing of Hartree-Fock and local density-functional theories, J. Chem. Phys. 98 (1993), pp. 1372–1377.
- [28] H.R. Karfunkel and R.J. Gdanitz, Ab initio prediction of possible crystal-structures for general organic-molecules, J. Comp. Chem. 13 (1992), pp. 1171–1183.
- [29] H.R. Karfunkel, F.J.J. Leusen, and R.J. Gdanitz, The ab initio prediction of yet unknown molecular crystal structures by solving the crystal packing problem, J. Comput Aided Mater. Design 1 (1994), pp. 177–185.
- [30] R.J. Gdanitz, *Prediction of molecular-crystal structures by Monte-Carlo simulated annealing without reference to diffraction data*, Chem. Phys. Lett. 190 (1992), pp. 391–396.
- [31] P. Verwer and F.J.J. Leusen, *Computer simulation to predict possible crystal polymorphs*, Rev. Comput. Chem. (1998), pp. 327–365.
- [32] A.L. Spek, Single-crystal structure validation with the program *PLATON*, J. Appl. Cryst. 36 (2003), pp. 7–13.
- [33] P.G. Karamertzanis, and S. L. Price, *Energy minimization of crystal structures containing flexible molecules*, J. Chem. Theory Comput. 2 (2006), pp. 1184–1199.
- [34] G.R. Desiraju, The Crystal as a Supramolecular Entity, John Wiley & Sons Ltd, Chichester, 1996.